

## INTEGRAL OPERATOR AND OSCILLATION OF SECOND ORDER ELLIPTIC EQUATIONS

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**Abstract.** By using integral operator, some oscillation criteria for second order elliptic differential equation

$$\sum_{i,j=1}^d D_i[A_{ij}(x)D_j y] + q(x)f(y) = 0, \quad x \in \Omega, \quad (E)$$

are established. The results obtained here can be regarded as the extension of the well-known Kamenev theorem to Eq.(E).

### 1. Introduction and Preliminaries

We are concerned with the oscillatory behavior of second order elliptic differential equation of the form

$$\sum_{i,j=1}^d D_i[A_{ij}(x)D_j y] + q(x)f(y) = 0, \quad (1)$$

where  $x = (x_1, \dots, x_d) \in \Omega(a) \subseteq \mathbb{R}^d$ ,  $d \geq 2$ ,  $|x| = [\sum_{i=1}^d x_i^2]^{1/2}$ ,  $D_i y = \partial y / \partial x_i$  for all  $i$ ,  $\Omega(a) = \{x \in \mathbb{R}^d : |x| \geq a\}$  for some  $a > 0$ .

Throughout this paper, we assume that the following conditions holds:

(A<sub>1</sub>)  $A = (A_{ij})$  is a real symmetric positive definite matrix function with  $A_{ij} \in C_{loc}^{1+\nu}(\Omega(a), \mathbb{R})$  for all  $i, j$ , and  $\nu \in (0, 1)$ .

Denote by  $u_{\max}(x)$  the largest eigenvalue of the matrix  $A$ . We suppose that there exists a function  $u \in C([a, \infty), \mathbb{R}^+)$  such that

$$u(r) \geq \max_{|x|=r} u_{\max}(x), \quad \text{for } r \geq a;$$

(A<sub>2</sub>)  $q \in C_{loc}^{\nu}(\Omega(a), \mathbb{R})$ ,  $\nu \in (0, 1)$ , and  $q(x)$  does not eventually vanish;

(A<sub>3</sub>)  $f \in C(\mathbb{R}, \mathbb{R}) \cup C^1(\mathbb{R} - \{0\}, \mathbb{R})$ ,  $yf(y) > 0$  whenever  $y \neq 0$ , and  $f'(y) \geq k > 0$  for all  $y \neq 0$ .

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As usual, a function  $y \in C_{loc}^{2+\nu}(\Omega(a), \mathbb{R})$ ,  $\nu \in (0, 1)$ , is called a solution of Eq.(1) if  $y(x)$  satisfies Eq.(1) for all  $x \in \Omega(a)$ . We restrict our attention only the nontrivial solution of Eq.(1), i.e., to the solution  $y(x)$  such that  $\sup\{|y(x)| : x \in \Omega(b)\} > 0$  for every  $b \geq a$ . Regarding the question of existence of solution of Eq.(1) we refer the reader to the monograph [2]. A nontrivial solution  $y(x)$  of Eq.(1) is said to be oscillatory in  $\Omega(a)$  if the set  $\{x \in \Omega(a) : y(x) = 0\}$  is unbounded, otherwise it is said to be nonoscillatory. Eq.(1) called oscillatory if all its nontrivial solutions are oscillatory.

Investigation of oscillation of Eq.(1) with variable coefficient  $q(x)$  was initiated by Noussair and Swanson [5], who first extended the well-known Fite-Wintner Theorem (see Fite [1] and Wintner [10]) to Eq.(1). As an excellent survey paper, the reader is to recommended to Swanson [6]. Recently, Xu [7], [8] and Zhang et al. [9] have discussed the oscillation property of Eq.(1) under the assumption that the function  $q(x)$  is an ‘‘integrally small’’ coefficient, and have shown that most of Kamenev’s results in [3] do hold equally well in the case for Eq.(1). However, as far as the authors know there are few results by using integral averaging techniques [4], even through the function  $q(x)$  is nonnegative. Motivated by this fact, we intend here to establish Kamenev’s integral theorem for oscillation of Eq.(1) based on the integral operator [11]. We are especially interested in case where  $q$  may be take on negative values for arbitrarily large  $|x|$ . Our methodology is somewhat different from that of previous authors, we believe that our approach is simpler and also provides a more unified for the study of Kamenev-type oscillation theorems.

Now, we introduce the general mean and the integral operator.

Following Wong [11], let  $D = \{(r, s) : r \geq s \geq a\} \in \mathbb{R}^2$ , and  $D_0 = \{(r, s) : r > s \geq a\} \in \mathbb{R}^2$ . Consider the kernel function  $H(r, s) \in C^1(D, \mathbb{R})$  such that the following conditions are satisfied:

(H<sub>1</sub>)  $H(r, r) = 0$  for  $r \geq a$ ;  $H(r, s) > 0$  on  $(r, s) \in D_0$ ;

(H<sub>2</sub>) For each  $s \geq a$ ,  $\lim_{r \rightarrow \infty} H(r, s) = \infty$ , and there exists a positive constant  $k_0$  such that

$$\lim_{r \rightarrow \infty} \frac{H(r, s)}{H(r, a)} = k_0, \quad \text{for all } s \geq a;$$

(H<sub>3</sub>)  $\frac{\partial H}{\partial s}(r, s) \leq 0$ ,  $-\frac{\partial H}{\partial s}(r, s) = h(r, s)H(r, s)$ , and  $\frac{\partial h}{\partial r}(r, s) \leq 0$ , for  $(r, s) \in D_0$ .

Let  $\rho \in C^1([a, \infty), \mathbb{R}^+)$ , we define an integral operator  $\Pi_\tau^\rho$ , which is defined in [11], in terms of  $H(r, s)$  and  $\rho(s)$  as

$$\Pi_\tau^\rho(\Theta; r) = \int_\tau^r H(r, s)\Theta(s)\rho(s)ds, \quad \text{for } r > \tau \geq a, \quad (2)$$

where  $\Theta \in C([a, \infty), \mathbb{R})$ .

For notational simplicity, for arbitrary given functions  $\Phi \in C^1([a, \infty), \mathbb{R}^+)$  and  $(u\phi) \in C^1([a, \infty), \mathbb{R})$ , we define

$$Q(r) = \Phi(r) \left\{ \int_{S_r} q(x) d\sigma + \frac{k}{\omega} u(r)\phi^2(r)r^{1-d} - [u(r)\phi(r)]' \right\},$$

$$p(r) = - \left[ \frac{\Phi'(r)}{\Phi(r)} + \frac{2k}{\omega} \phi(r)r^{1-d} \right], \quad g(r) = \frac{kr^{1-d}}{\omega u(r)\Phi(r)}, \quad (3)$$

where  $S_r = \{x \in \mathbb{R}^d : |x| = r\}$  for  $r > a$ ,  $d\sigma$  denotes the spherical integral element in  $\mathbb{R}^d$ ,  $\omega$  denotes the surface area of unit sphere in  $\mathbb{R}^d$ , i.e.,  $\omega = 2\pi^{\frac{d}{2}}/\Gamma(\frac{d}{2})$ .

**Lemma 1.**[11] *Let  $H(r, s) \in C^1(D, \mathbb{R})$  satisfying  $(H_1)$ – $(H_3)$  and  $\rho(s) \in C([a, \infty), \mathbb{R})$ .*

*Then*

- (i)  $\lim_{r \rightarrow \infty} \frac{1}{H(r, a)} \int_a^r H(r, s)\rho(s) ds = \infty$ , if  $\lim_{r \rightarrow \infty} \int_a^r \rho(s) ds = \infty$ ;
- (ii)  $\frac{1}{H(r, a)} \int_a^r H(r, s)\rho(s) ds$  is nondecreasing in  $r$ , if  $\rho(s) \geq 0$  on  $[a, \infty)$ .

## 2. Main Results

In this section, we establish Kamenev's integral oscillation criteria for Eq.(1) based on the integral operator  $\Pi_r^\rho$ . For the sake of simplicity, we always assume that the kernel function  $H(r, s)$  satisfies condition  $(H_1)$  –  $(H_3)$ , the integral operator  $\Pi_r^\rho$  and functions  $Q, p, g$  defined by (2) and (3), respectively.

**Theorem 1.** *Assume that there exist functions  $\Phi \in C^1([a, \infty), \mathbb{R}^+)$  and  $(u\phi) \in C^1([a, \infty), \mathbb{R})$  such that*

$$\limsup_{r \rightarrow \infty} \frac{1}{H(r, a)} \Pi_a^\tau \left( Q - \frac{1}{4g} \left[ h + p - \frac{\rho'}{\rho} \right]^2 \right) = \infty. \quad (4)$$

*Then Eq.(1) is oscillatory.*

**Proof.** Let  $y(x)$  be a nonoscillatory solution of Eq.(1). Without loss of generality, we may assume that  $y(x) > 0$  for all  $x \in \Omega(a)$ . Define

$$W(x) = \frac{1}{f(y)} (A^\nabla y)(x), \quad (5)$$

where  $\nabla y = (D_1 y, D_2 y, \dots, D_d y)^T$  denotes the gradient of  $y(x)$ . Differentiation (5) and using Eq.(1), we obtain that

$$\operatorname{div} W(x) = -q(x) - f'(y)(W^T A^{-1} W)(x).$$

Put

$$Z(r) = \Phi(r) \left[ \int_{S_r} W(x) \cdot \mu(x) d\sigma + u(r)\phi(r) \right], \quad (6)$$

where  $\mu(x) = x/|x|$ , ( $x \neq 0$ ), denotes the outward unit normal. By Green's formula in (6), we have

$$\begin{aligned} Z'(r) &= \frac{\Phi'(r)}{\Phi(r)}Z(r) + \Phi(r) \left\{ \int_{S_r} \operatorname{div}W(x) d\sigma + [u(r)\phi(r)]' \right\} \\ &= \frac{\Phi'(r)}{\Phi(r)}Z(r) - \Phi(r) \left\{ \int_{S_r} q(x)d\sigma + \int_{S_r} f'(y)(W^T A^{-1}W)(x) d\sigma - [u(r)\phi(r)]' \right\}. \end{aligned} \quad (7)$$

In view of  $(A_1)$ ,

$$(W^T A^{-1}W)(x) \geq u_{\max}^{-1}(x)|W(x)|^2.$$

The Schwartz inequality gives

$$\int_{S_r} |W(x)|^2 d\sigma \geq \frac{r^{1-d}}{\omega} \left[ \int_{S_r} W(x) \cdot \mu(x) d\sigma \right]^2.$$

Thus, by (6) and (7), we obtain

$$\begin{aligned} Z'(r) &\leq \frac{\Phi'(r)}{\Phi(r)}Z(r) - \Phi(r) \left\{ \int_{S_r} q(x) d\sigma + \frac{kr^{1-d}}{\omega u(r)} \left[ \int_{S_r} W(x) \cdot \mu(x) d\sigma \right]^2 - [u(r)\phi(r)]' \right\} \\ &= \frac{\Phi'(r)}{\Phi(r)}Z(r) - \Phi(r) \left\{ \int_{S_r} q(x) d\sigma + \frac{kr^{1-d}}{\omega u(r)} \left[ \frac{Z(r)}{\Phi(r)} - u(r)\phi(r) \right]^2 - [u(r)\phi(r)]' \right\} \\ &= -\Phi(r) \left\{ \int_{S_r} q(x) d\sigma + \frac{k}{\omega} u(r)\phi^2(r)r^{1-d} - [u(r)\phi(r)]' \right\} \\ &\quad + \left[ \frac{\Phi'(r)}{\Phi(r)} + \frac{2k}{\omega} \phi(r)r^{1-d} \right] Z(r) - \frac{kr^{1-d}}{\omega u(r)\Phi(r)} Z^2(r). \end{aligned} \quad (8)$$

Denote  $V(r) = Z(r) + \frac{p(r)}{2g(r)}$ , (8) can be rewritten as

$$Z'(r) + g(r)V^2(r) + Q(r) - \frac{1}{4} \frac{p^2(r)}{g(r)} \leq 0. \quad (9)$$

Applying the integral operator  $\Pi_b^\rho$ , ( $b \geq a$ ), to (9), we obtain

$$\Pi_b^\rho(gV^2) + \Pi_\tau^\rho \left( \left[ h - \frac{\rho'}{\rho} \right] Z \right) + \Pi_\tau^\rho \left( Q - \frac{p^2}{4g} \right) \leq H(r, b)\rho(b)Z(b). \quad (10)$$

Completing squares of  $V$  in (10) yields

$$\Pi_b^\rho \left( g \left\{ V + \frac{1}{2g} \left[ h - \frac{\rho'}{\rho} \right] \right\}^2 \right) + \Pi_b^\rho \left( Q - \frac{1}{4g} \left[ h + p - \frac{\rho'}{\rho} \right]^2 \right) \leq H(r, b)\rho(b)Z(b). \quad (11)$$

Let  $b = a$  and divide (11) through by  $H(r, a)$ . Note the first term is nonnegative, thus

$$\frac{1}{H(r, a)} \Pi_a^\tau \left( Q - \frac{1}{4g} \left[ h + p - \frac{\rho'}{\rho} \right]^2 \right) \leq \rho(a)Z(a). \quad (12)$$

Take linsup in (12) as  $r \rightarrow \infty$ . Condition (4) gives a desired contradiction in (12). This proves Theorem 1.

**Remark 1.** If we choose the function  $\phi(s) \equiv 0$ , then Theorem 1 contains Theorem 4 in [5].

**Theorem 2.** Let  $\Phi$  and  $\phi$  be as Theorem 1. Assume that there exist functions  $\varphi_i \in C([a, \infty), \mathbb{R})$ , ( $i = 1, 2$ ), such that for  $\tau \geq a$ ,

$$\lim_{r \rightarrow \infty} \frac{1}{H(r, a)} \Pi_\tau^\rho \left( \frac{1}{g} \left[ h + p - \frac{\rho'}{\rho} \right]^2 \right) \leq \varphi_1(\tau), \quad (13)$$

and

$$\limsup_{r \rightarrow \infty} \frac{1}{H(r, a)} \Pi_\tau^\rho(Q) \geq \varphi_2(\tau), \quad (14)$$

where  $\varphi_1$  and  $\varphi_2$  satisfy

$$\lim_{r \rightarrow \infty} \frac{1}{H(r, a)} \Pi_a^\rho \left( g\rho^{-2} \left[ \varphi_2 - \frac{1}{4}\varphi_1 \right]_+^2 \right) = \infty, \quad (15)$$

and  $[\varphi(r)]_+ = \max\{\varphi(r), 0\}$ . Then Eq.(1) is oscillatory.

**Proof.** Proceeding as in the proof Theorem 1, we see (11) holds for all  $r \geq a$ . Divide (11) by  $H(r, a)$  and drop the nonnegative first term and obtain

$$\frac{1}{H(r, a)} \Pi_\tau^\rho(Q) - \frac{1}{4} \frac{1}{H(r, a)} \Pi_\tau^\rho \left( \frac{1}{g} \left[ h + p - \frac{\rho'}{\rho} \right]^2 \right) \leq \frac{H(r, b)}{H(r, a)} \rho(\tau) Z(\tau).$$

Take linsup in above inequality as  $r \rightarrow \infty$  and note from (13), (14) and  $(H_2)$  that

$$\varphi_2(\tau) - \frac{1}{4}\varphi_1(\tau) \leq k_0 \rho(\tau) Z(\tau).$$

From which it follows that

$$k_0^{-2} \rho^{-2}(\tau) \left[ \varphi_2(\tau) - \frac{1}{4}\varphi_1(\tau) \right]_+^2 \leq Z^2(\tau). \quad (16)$$

Then, it follows from (15) and (16) that

$$\lim_{r \rightarrow \infty} \frac{1}{H(r, a)} \Pi_a^\rho(gZ^2) \geq \lim_{r \rightarrow \infty} \frac{k_0^{-2}}{H(r, a)} \Pi_a^\rho \left( g\rho^{-2} \left[ \varphi_2 - \frac{1}{4}\varphi_1 \right]_+^2 \right) = \infty. \quad (17)$$

Next, we shall show that (17) is not possible.

For (11), set  $b = a$ . Dividing (11) through  $H(r, a)$ , and noting (13), (14), we have

$$\frac{1}{H(r, a)} \Pi_a^\rho \left( g \left\{ V + \frac{1}{2g} \left[ h - \frac{\rho'}{\rho} \right] \right\}^2 \right) \leq \rho(a) Z(a) + \left[ \varphi_2(b) - \frac{1}{4}\varphi_1(b) \right]. \quad (18)$$

We note that by Lemma 1 (ii), (18) and (13), the following limits exist and finite, that is

$$\begin{aligned} \lim_{r \rightarrow \infty} \frac{1}{H(r, a)} \Pi_a^\rho \left( g \left\{ V + \frac{1}{2g} \left[ h - \frac{\rho'}{\rho} \right] \right\}^2 \right) < \infty, \\ \lim_{r \rightarrow \infty} \frac{1}{H(r, a)} \Pi_a^\rho \left( \frac{1}{g} \left[ h + p - \frac{\rho'}{\rho} \right]^2 \right) < \infty. \end{aligned} \quad (19)$$

Noting that

$$Z(r) = V(r) - \frac{p(r)}{2g(r)} = \left( V + \frac{1}{2g} \left[ h - \frac{\rho'}{\rho} \right] \right) - \frac{1}{2g} \left( h + p - \frac{\rho'}{\rho} \right).$$

Thus

$$g(s)Z^2(r) \leq 2 \left\{ g(s) \left( V + \frac{1}{2g} \left[ h - \frac{\rho'}{\rho} \right] \right)^2 + \frac{1}{4g} \left( h + p - \frac{\rho'}{\rho} \right)^2 \right\}.$$

Consequently, by (19), we obtain

$$\lim_{r \rightarrow \infty} \frac{1}{H(r, a)} \Pi_a^\rho (gZ^2) < \infty,$$

which contradicts (17). Therefore, the proof is complete.

**Remark 2.** Note that Theorem 1 and Theorem 2 do not require conditions  $\int^\infty \frac{r^{1-d}}{u(r)} dr = \infty$ , or  $\int_{\Omega(b)} q(x) dx$ , ( $b \geq a$ ), convergent in [5], [7-9].

In same way as it was done in [11], with an appropriate choice of the functions  $H$  and  $\rho$ , we can derive various oscillation criteria for Eq.(1) from Theorems 1 and 2. For example, define

$$H(r, s) = (r - s)^\alpha, \quad (r, s) \in D, \quad \alpha > 1.$$

**Corollary 1.** Let  $\alpha > 1$ . Suppose that the function  $g$  is nondecreasing, and

$$\limsup_{r \rightarrow \infty} r^{-\alpha} \int_a^r (r - s)^\alpha \left\{ Q(s) - \frac{p^2(s)}{4g(s)} - \frac{1}{2} \left( \frac{p(s)}{g(s)} \right)' \right\} ds = \infty. \quad (20)$$

Then Eq.(1) is oscillatory.

**Proof.** Let  $\rho(r) = 1$  in Theorem 1, observe that

$$[p(s) + \alpha(r - s)^{-1}]^2 = p^2(s) + 2\alpha(r - s)^{-1}p(s) + \alpha^2(r - s)^{-2}.$$

By the Bonnet Theorem, for a fixed  $r \geq a$  and some  $\xi \in [a, r]$ , we have

$$\int_a^r \frac{(r - s)^{\alpha-2}}{g(s)} ds = \frac{1}{g(a)} \int_a^\xi (r - s)^{\alpha-2} ds,$$

hence

$$\lim_{r \rightarrow \infty} \frac{1}{r^\alpha} \int_a^r \frac{(r-s)^{\alpha-2}}{g(s)} ds = 0.$$

Using integrating by parts, we get

$$\alpha \int_a^r (r-s)^{\alpha-1} \frac{p(s)}{g(s)} ds = (r-a)^\alpha \frac{p(a)}{g(a)} + \int_a^r (r-s)^\alpha \left( \frac{p(s)}{g(s)} \right)' ds.$$

Thus, Corollary 1 from Theorem 1.

Now, let  $\rho(r) = \exp(\int_a^r p(s) ds)$  in condition (4), we have

**Corollary 2.** *Let  $\alpha > 1$ . If*

$$\limsup_{r \rightarrow \infty} r^{-\alpha} \int_a^r (r-s)^\alpha Q(s) \exp \left( \int_a^s p(s) d\tau \right) ds = \infty, \quad (21)$$

and

$$\lim_{r \rightarrow \infty} r^{-\alpha} \int_a^r (r-s)^{\alpha-2} [g(s)]^{-1} \exp \left[ \int_a^s p(\tau) d\tau \right] ds < \infty. \quad (22)$$

Then Eq.(1) is oscillatory.

**Remark 3.** With the same choice of the functions  $H(r, s)$  and  $\rho(s)$  as corollary 1 and Corollary 2, more general oscillation criteria for Eq.(1) can be obtained by Theorem 2 similarly. Here we omit the details.

**Example 1.** Consider the linear elliptic equation of second order

$$\Delta y + (|x|^{\sigma-1} \sin |x|)y = 0, \quad (23)$$

where  $x \in \Omega(1)$ ,  $d = 2$ ,  $\sigma > 0$ ,  $u(r) = 1$ ,  $q(x) = |x|^{\sigma-1} \sin |x|$ ,  $f'(y) = 1$  for all  $|x| > 1$ . Choose  $\Phi(r) = 1$ ,  $\varphi(r) = 0$ , then

$$Q(r) = 2\pi r^\sigma \sin r, \quad g(r) = (2\pi r)^{-1}, \quad p(s) = 0.$$

For  $\sigma > 0$ , it is easy to verify that

$$\limsup_{r \rightarrow \infty} \frac{1}{r^2} \int_1^r (r-s)^2 Q(s) ds = \limsup_{r \rightarrow \infty} \frac{1}{r^2} \int_1^r (r-s)^2 s^\sigma \sin s ds = \infty,$$

and

$$\lim_{r \rightarrow \infty} \frac{1}{r^2} \int_1^r \frac{1}{g(s)} ds = \pi.$$

Hence, Eq.(23) is oscillatory by Corollary 2.

**Example 2.** Consider the nonlinear elliptic equation of second order

$$\Delta y + (|x|^v \cos |x| + \sin |x|)(y + y^3) = 0, \quad (24)$$

where  $x \in \Omega(1)$ ,  $d = 2$ ,  $v$  is a constant with  $2v > -1$ ,  $u(r) = 1$ ,  $q(x) = |x|^v \cos |x| + \sin |x|$ , and  $f(y) = y + y^3$ . Let  $\Phi(r) = r^{-1}$ ,  $\varphi(r) = 0$ , then

$$g(r) = (2\pi)^{-1}, \quad p(r) = r^{-1}, \quad Q(r) = 2\pi(r^v \cos r + \sin r), \quad f'(y) = 1 + y^2.$$

Taking  $H(r, s) = (r - s)^2$  for  $(r, s) \in D$  and  $\rho(r) = 1$ , we have, for all  $\tau \geq 1$ ,

$$\limsup_{r \rightarrow \infty} \frac{1}{r^2} \int_{\tau}^r (r - s)^2 Q(s) ds \geq -2\pi\tau^v \cos \tau - \varepsilon,$$

and

$$\lim_{r \rightarrow \infty} \frac{1}{r^2} \int_{\tau}^r \frac{1}{g(s)} (r - s)^2 [2(r - s)^{-1} + s^{-1}]^2 ds = \frac{2\pi}{\tau},$$

where  $\varepsilon$  is a positive constant.

Now, set  $\varphi_1(\tau) = (2\pi\tau)^{-1}$  and  $\varphi_2(\tau) = -2\pi\tau^v \cos \tau - \varepsilon$ . Then, there exists an integer  $N$  such that  $(2N + 1)\pi - \frac{\pi}{4} > 1$  and if  $n \geq N$ ,

$$(2n + 1)\pi - \frac{\pi}{4} \leq \tau \leq (2n + 1)\pi + \frac{\pi}{4}, \quad \varphi_2(\tau) - \frac{1}{4}\varphi_1(\tau) \geq \delta\tau^v,$$

where  $\delta$  is a small constant. Taking into account that  $2v > -1$ , we obtain

$$\begin{aligned} \lim_{r \rightarrow \infty} \int_a^r \left[ \varphi_2(\tau) - \frac{1}{4}\varphi_1(\tau) \right]_+^2 ds &\geq \sum_{n=N}^{\infty} \delta^2 \int_{(2n+1)\pi - \frac{\pi}{4}}^{(2n+1)\pi + \frac{\pi}{4}} s^{2v} ds \\ &\geq \sum_{n=N}^{\infty} \delta^2 \int_{(2n+1)\pi - \frac{\pi}{4}}^{(2n+1)\pi + \frac{\pi}{4}} s^{-1} ds = \infty. \end{aligned}$$

This implies that, by Lemma 1 (i),

$$\lim_{r \rightarrow \infty} \frac{1}{r^2} \int_a^r \frac{1}{2\pi} (r - s)^2 \left[ \varphi_2(\tau) - \frac{1}{4}\varphi_1(\tau) \right]_+^2 ds = \infty.$$

Accordingly, all conditions of Theorem 2 are satisfied, and hence, Eq.(24) is oscillatory.

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