



## THE FORMULA FOR THE REGULARIZED TRACE OF THE STURM-LIOUVILLE OPERATOR WITH A LOGARITHMIC POTENTIAL

KHABIR KABIROVICH ISHKIN AND LEISAN GABDULLOVNA VALIULLINA

**Abstract.** We have obtained a regularized trace formula for the Sturm-Liouville operator on a semi-axis with a logarithmic potential.

### 1. Introduction

Let the function  $q$  be defined on  $(0, +\infty)$ , real-valued, summable on any finite interval  $(0, b)$ ,  $b > 0$ , and

$$\lim_{x \rightarrow +\infty} q(x) = +\infty.$$

Further, let  $L$  be the operator generated in  $L^2(, +\infty)$  by the differential expression  $ly := -y'' + qy$  and the boundary condition  $y(0) = 0$  [1, Ch. V, § 18]. According to the well-known A. M. Molchanov theorem [2], operator  $L$  has a discrete spectrum. In the paper [3] an asymptotic equation for the spectrum of the operator  $L$  whose potential can grow arbitrarily slowly was obtained. This equation allows us to calculate the first few (up to a summable remainder) terms of the asymptotic series for the eigenvalues in the case

$$q = \underbrace{\log \dots \log x}_m, m \in \mathbb{N}, a = \text{const} > \begin{cases} e^{m-2}, & m \geq 2, \\ 0, & m = 1. \end{cases}$$

In particular, when  $q = \log(x + a)$

$$\lambda_k = s_k + O(k^{-1}(\log k)^{-3/2}), \tag{1}$$

$$s_k = \log(2\sqrt{\pi}k) - k^{-1} \left( \frac{a}{\pi} \sqrt{\log k} + k_0 - 1/4 + \frac{c_0}{\sqrt{\log k}} \right), \tag{2}$$

where  $c_0 = \frac{a}{2\pi} (1 + \log(2\sqrt{\pi}/a))$ . A parameter  $k_0$  is some positive integer (regularization defect) that ensures the convergence of the series

$$\sum_{k=2}^{\infty} (\lambda_k - s_k). \tag{3}$$

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The purpose of this paper is to find the value of  $k_0$  and to calculate the sum of the series (3), which is called the regularized trace of the operator  $L_a$  with the potential  $q = \log(x + a)$ .

The first  $L$  operator with the potential  $q(x) = \log^+ x := \max\{\log x, 0\}$  was considered in [4] as an example of Feynman–Kac theory application [5, § X.11] (in combination with Karamata’s Tauberian theorem [6, Ch. XII, § 7]) to the problem of finding the asymptotics of the function

$$N(\lambda) := \sum_{\lambda_k < \lambda} 1. \tag{4}$$

This result was summarized by K. Kh. Boymatov [7] on Sturm–Liouville operators with a matrix potential that allows growth of order  $\underbrace{\log \dots \log x}_m$  ( $m \in \mathbb{N}$ ).

Regularized traces of the form (3) in the case of power growth potentials are well studied (see [8] and references therein). From the formulas (1) – (2) it is clear that for any  $n$  operator  $L_a^{-n}$  is not a trace-class operator, therefore the classical method of zeta-functions [9, 10] is not applicable in this situation. On the other hand, due to the exponential growth of the function (4)  $\theta$ -function of the operator  $L_a$

$$\Theta_a(t) := \sum_{k=1}^{\infty} e^{-t\lambda_k}. \tag{5}$$

is determined only on the half-plane  $\Re t > c$  with some positive  $c$ . Therefore, the parabolic equation method [11] using asymptotics  $\Theta_a(t)$  as  $t \rightarrow +0$  seems inapplicable in this case. However, formulas (1) – (2) allow us to construct an analytic continuation of the function  $\Theta_a$  on the half-plane  $\Re t > 0$ . Using this fact, it is possible to find the sum of a series (3).

## 2. Asymptotics of the function $\Theta_a$

By virtue of formulas (1) – (2) and (5) the function  $\Theta_a$  holomorphic in the half-plane  $\Re t > 1$ .

**Theorem 1.** *The function  $\Theta_a$  admits a meromorphic (with a single pole at  $t = 1$ ) continuation to the half-plane  $\Re t > 0$  and as  $t \rightarrow +0$  the following asymptotic decomposition*

$$\Theta_a(t) \sim -\frac{a}{2\sqrt{\pi}} t^{-1/2} - (k_0 + 1/4) + \theta_0 t^{1/2} + \theta_1 t + O(t^{3/2}), \tag{6}$$

where

$$\begin{aligned} \theta_0 &= \frac{a(\log a - 1)}{2\sqrt{\pi}}, \\ \theta_1 &= \lambda_1 + \sigma(2) + \frac{a}{\pi} \int_1^{\infty} \{x\} (\log x)^{1/2} x^{-1} dx \\ &\quad + \left(k_0 - \frac{1}{4}\right)(\gamma - 1) - \left(k_0 + \frac{5}{4}\right) \log 2\sqrt{\pi} + \frac{1}{2} \log 2\pi \end{aligned}$$

$$+ \frac{a}{2\pi} \left( 1 + \log \frac{2\sqrt{\pi}}{a} \right) \int_1^\infty \{x\} (\log x)^{-1/2} x^{-1} dx,$$

$\gamma$  – Euler’s constant.

**Proof.** We transform the decomposition (1) – (2) to a form, which is more convenient for finding the asymptotics of the function  $\Theta_a$ :

$$\lambda_k = \log k + c_1 + c_2(\log k)^{1/2}k^{-1} + c_3k^{-1} + c_0(\log k)^{-1/2}k^{-1} + r_k, \tag{7}$$

where  $c_1 = \log 2\sqrt{\pi}$ ,  $c_2 = \frac{a}{\pi}$ ,  $c_3 = k_0 - \frac{1}{4}$ ,  $c_0 = \frac{a}{2\pi} (1 + \log(2\sqrt{\pi}/a))$ ,

$$r_k = O(k^{-1}(\log k)^{-3/2}), \quad k \rightarrow \infty. \tag{8}$$

Then

$$\begin{aligned} \Theta_a(t) &= \sum_{k=1}^\infty e^{-t\lambda_k} = e^{-t\lambda_1} + e^{-tc_1} \sum_{k=2}^\infty e^{-t(\log k + c_2(\log k)^{1/2} + c_3k^{-1} + c_0(\log k)^{-1/2}k^{-1} + r_k)} \\ &= e^{-t\lambda_1} + e^{-tc_1} \sum_{k=2}^\infty k^{-t} e^{-c_2t(\log k)^{1/2}k^{-1}} e^{-c_3tk^{-1}} e^{-c_0t(\log k)^{-1/2}k^{-1}} e^{-tr_k}. \end{aligned} \tag{9}$$

Putting each exponent according to the Taylor formula, we get

$$\begin{aligned} e^{-c_3tk^{-1}} &= 1 - c_3tk^{-1} + O(t^2k^{-2}); \\ e^{-c_0t(\log k)^{-1/2}k^{-1}} &= 1 - c_0t(\log k)^{-1/2}k^{-1} + O(t^2(\log k)^{-1}k^{-2}); \\ e^{-c_2t(\log k)^{1/2}k^{-1}} &= 1 - c_2t(\log k)^{1/2}k^{-1} + O(t^2(\log k)k^{-2}); \\ e^{-tr_k} &= 1 - tr_k + O(t^2r_k^2). \end{aligned}$$

Substituting these expansions into (9), we have

$$\begin{aligned} \Theta_a(t) &= e^{-t\lambda_1} + e^{-tc_1} \{ \zeta(t) - 1 - t [c_2\varphi_1(t) + c_3(\zeta(t+1) - 1) \\ &\quad + c_0\varphi_2(t) + \varphi_3(t)] \} + \varphi_4(t), \end{aligned} \tag{10}$$

where

$$\begin{aligned} \varphi_1(t) &= \sum_{k=2}^\infty (\log k)^{1/2}k^{-1-t}, & \varphi_2(t) &= \sum_{k=2}^\infty (\log k)^{-1/2}k^{-1-t}, \\ \varphi_3(t) &= \sum_{k=2}^\infty k^{-t}r_k, & \varphi_4(t) &= e^{-tc_1} \sum_{k=2}^\infty k^{-t}R_k(t), \\ R_k(t) &= O(t^2(\log k)k^{-2}). \end{aligned}$$

Therefore, the function  $\Theta_a$  admits a meromorphic extension to the half-plane  $\Re t > 0$  with a single pole at  $t = 1$  (because of the term  $\zeta(t)$ ). Find the asymptotics  $\Theta_a(t)$  as  $t \rightarrow +0$ .

Let  $\rho(x) = x - \{x\}$ , where  $\{x\}$  is fractional part of the number  $x$ . Then  $\forall 0 < \varepsilon < 1$

$$\begin{aligned} \varphi_1(t) &= \int_{1+\varepsilon}^{\infty} (\log x)^{1/2} x^{-1-t} d\rho(x) = \int_{1+\varepsilon}^{\infty} (\log x)^{1/2} x^{-1-t} dx \\ &\quad + \varepsilon (\log(1 + \varepsilon))^{1/2} (1 + \varepsilon)^{-1-t} + \int_{1+\varepsilon}^{\infty} \{x\} [(\log x)^{1/2} x^{-1-t}]' dx, \end{aligned}$$

whence as  $\varepsilon \rightarrow +0$  we get

$$\varphi_1(t) = \int_1^{\infty} (\log x)^{1/2} x^{-1-t} dt + \int_1^{\infty} \{x\} ((\log x)^{1/2} x^{-1-t})' dx =: \varphi_{11}(t) + \varphi_{12}(t).$$

Direct calculation gives

$$\varphi_{11}(t) = \frac{\sqrt{\pi}}{2} t^{-3/2}.$$

Further,

$$\varphi_{12}(t) = \int_1^{\infty} \{x\} \left[ -(1+t)(\log x)^{1/2} + \frac{1}{2}(\log x)^{-1/2} \right] x^{-2-t} dx.$$

Denote the integrand by  $f(x, t)$ . Then  $\forall k$

$$\frac{\partial^k}{\partial t^k} f(x, t) = O\left((\log x)^{-1/2+k} x^{-2}\right), \quad x \geq 1,$$

uniformly in  $\Re t > 0$ , therefore  $\varphi_{12}$  is holomorphic at zero, so

$$\varphi_{1,2}(t) = \varphi_{1,2}(0) + O(t), \quad t \rightarrow 0.$$

Therefore,

$$\varphi_1(t) = \frac{\sqrt{\pi}}{2} t^{-3/2} + \int_1^{\infty} \{x\} ((\log x)^{1/2} x^{-1})' dx + O(t), \quad t \rightarrow +0. \tag{11}$$

Similarly, it is proved that

$$\varphi_2(t) = \sqrt{\pi} t^{-1/2} + \int_1^{\infty} \{x\} ((\log x)^{-1/2} x^{-1})' dx. \tag{12}$$

Let us introduce the function

$$\sigma(x) = \sum_{k \geq x} r_k.$$

From (1) – (2) it follows that

$$\sigma(x) = O((\log x)^{-1/2}), \quad x \rightarrow \infty.$$

Then

$$\varphi_3(t) = - \int_2^{\infty} x^{-t} d\sigma(x) = \frac{\sigma(2)}{2^t} - t \int_2^{\infty} x^{-t-1} \sigma(x) dx = \sigma(2) + O(t^{1/2}), \quad t \rightarrow +0. \tag{13}$$

Further, due to the well-known properties of  $\zeta$ -functions (see, for example, [12, Ch. II, 1<sup>0</sup>])

$$\zeta(t) = -\frac{1}{2} - \frac{1}{2} \log(2\pi)t + O(t^2), \tag{14}$$

$$\zeta(t+1) = \frac{1}{t} + \gamma + O(t). \tag{15}$$

Now substituting (11) – (15) into (10), after simple calculations we get (6). The theorem is proved.

### 3. Operator $e^{-tL_a}$ and its trace

#### 3.1. Estimation of the resolvent kernel for the operator $L_a$

We introduce the notation. Let be  $\Omega_a = \mathbb{C} \setminus [\log a, +\infty)$ ,  $\Omega_{r\sigma} = \{\sigma \leq \arg(\lambda - \log a - r) \leq 2\pi - \sigma\}$ ,  $r > 0, 0 < \sigma < \pi$ ,  $\varphi(x, \lambda)$  and  $\psi(x, \lambda)$  — solutions of the equation

$$-y'' + q_a y = \lambda y \tag{16}$$

that satisfy the following conditions:

$$\begin{aligned} \varphi(0, \lambda) = 0, \quad \frac{\partial \varphi}{\partial x}(0, \lambda) = 1, \\ \psi(x, \lambda) \sim (q_a(x) - \lambda)^{-1/4} \exp\left(-\int_0^x (q_a(t) - \lambda)^{1/2} dt\right), \quad x \rightarrow +\infty. \end{aligned} \tag{17}$$

Hereinafter, the expression  $z^{1/n}$  will mean that branch of the root  $\sqrt[n]{z}$ , which is positive for positive  $z$ . Since for each fixed  $\lambda \notin [q_a(b), +\infty)$  the function

$$\alpha(x, \lambda) = \frac{1}{8} \frac{q_a''(x)}{(q_a(x) - \lambda)^{3/2}} - \frac{5}{32} \frac{q_a'^2(x)}{(q_a(x) - \lambda)^{5/2}}, \tag{18}$$

is summable on  $(b, +\infty)$ , there exists a unique solution  $\psi$ , satisfying (17), [13, Ch. II, § 6].

The resolvent kernel for the operator  $L_a$  has the form

$$G(x, y, \lambda) = \frac{1}{\psi(0, \lambda)} \begin{cases} \varphi(x, \lambda)\psi(t, \lambda), & x < t < \infty, \\ \psi(x, \lambda)\varphi(t, \lambda), & 0 < t < x. \end{cases} \tag{19}$$

Therefore,

$$G(x, x, \lambda) = \frac{1}{\psi(0, \lambda)} \varphi(x, \lambda)\psi(x, \lambda). \tag{20}$$

**Lemma 1.** For each  $\varepsilon > 0$  there is a constant  $r_\varepsilon > 0$  such that if  $\varepsilon \leq \sigma < \pi, r \geq r_\varepsilon$ , then for all  $x \geq 0$  and  $\lambda \in \Omega_{r\sigma}$  function  $\psi(x, \lambda)$  permits the representation

$$\psi(x, \lambda) \sim (q_a(x) - \lambda)^{-1/4} \exp\left(-\int_0^x (q_a(t) - \lambda)^{1/2} dt\right) \times$$

$$\times \left[ 1 + \int_x^{+\infty} \left( \alpha + \frac{\alpha'}{2\sqrt{q-\lambda}} + R \right) dt \right], \tag{21}$$

where  $\alpha$  is defined by the formula (18) and

$$\sup_{x \geq 0, \lambda \in \Omega_{r\sigma}} |R(x, \lambda)(x+a)^3(q-\lambda)^{5/2}| < \infty. \tag{22}$$

**Proof.** Substituting

$$\psi(x, \lambda) = (q_a(x) - \lambda)^{-1/4} \exp \left( - \int_0^x (q_a(t) - \lambda)^{1/2} dt + \int_{x_0}^x (-\alpha + \beta) dt \right), \tag{23}$$

we get the equation for  $\beta$

$$\beta' - 2(\beta_0 + \beta_1 + \alpha)\beta + 2\beta_1\alpha + \alpha^2 - \alpha' + \beta^2 = 0,$$

where

$$\beta_0 = \sqrt{q_a - \lambda}, \beta_1 = \frac{1}{4} \frac{q'_a}{q_a - \lambda}.$$

The method of variation of constants leads to the equation

$$\beta = \gamma + \int_x^{+\infty} e^{-2 \int_x^t (\beta_0 + \alpha) dt} (q_a(x) - \lambda)^{1/2} (q_a(t) - \lambda)^{-1/2} \beta^2 dt, \tag{24}$$

where

$$\gamma = \int_x^{+\infty} e^{-2 \int_x^t (\beta_0 + \alpha) dt} (q_a(x) - \lambda)^{1/2} (q_a(t) - \lambda)^{-1/2} (2\beta_1\alpha - \alpha' + \alpha^2) dt.$$

By direct calculations, it is easy to verify that for every  $r > 0, 0 < \sigma < \pi$

$$|q_a(t) - \lambda| > \sin \sigma |q_a(x) - \lambda| \quad \forall t \geq x, \lambda \in \Omega_{r\sigma}.$$

Integrating in parts and taking into account the last inequality, we get

$$\gamma = - \frac{\alpha'}{2\sqrt{q_a - \lambda}} + O((x+a)^{-3}(q_a(x) - \lambda)^{-5/2})$$

uniformly in  $x \geq 0$  and  $\lambda \in \Omega_{r\sigma}$  for all  $r > 0, 0 < \sigma < \pi$ .

Replacing

$$R = (x+a)^3(q_a(x) - \lambda)^{5/2} \left( \beta + \frac{\alpha'}{2\sqrt{q_a - \lambda}} \right) \tag{25}$$

converts the equation (24) to the form

$$R(x, \lambda) = \tilde{\gamma}(x, \lambda) + \int_x^{+\infty} K(x, t, \lambda) R^2(t, \lambda) dt, \tag{26}$$

where functions  $\tilde{\gamma}(x, \lambda)$  and  $K(x, t, \lambda)$  are continuous on  $[0, +\infty) \times \Omega_a$  and  $[0, +\infty)^2 \times \Omega_a$  respectively. Moreover,

$$\tilde{\gamma}(x, \lambda) = O(1), K(x, t, \lambda) = O((x + a)^{-3}(q_a(x) - \lambda)^{-5/2})$$

uniformly in  $x \geq 0$  and  $\lambda \in \Omega_{r\sigma}$  (for all  $r > 0, 0 < \sigma < \pi$ ). According to the second estimate for any  $\varepsilon > 0$  there is a sufficiently large  $r_\varepsilon > 0$  such that for all  $r \geq r_\varepsilon$  and  $\varepsilon \leq \sigma < \pi$  the integral operator on the right side of (26) is a contraction in a Banach space  $C([0, +\infty) \times \Omega_{r\sigma})$ . The method of successive approximations shows that for all  $r \geq r_\varepsilon$  and  $\varepsilon \leq \sigma < \pi$ , the equation (26) has a unique solution  $R \in C([0, +\infty) \times \Omega_{r\sigma})$ . Hence, according to the equalities (25) and (23), for  $x_0 = +\infty$ , the representation (21) with the estimate (22) follows. The lemma is proved.

According to the proof of the lemma it follows that if  $\varepsilon \leq \sigma < \pi, r \geq r_\varepsilon$  then the function  $\psi(x, \lambda)$  for every  $\lambda \in \Omega_{r\sigma}$  does not have zeros on  $[0, +\infty)$ . Therefore, the solution  $\varphi(x, \lambda)$  can be accepted in the form

$$\varphi(x, \lambda) = \psi(0, \lambda)\psi(x, \lambda) \int_0^x \psi^{-2}(t, \lambda) dt. \tag{27}$$

Substituting this expression into (20), we get

$$G(x, x, \lambda) = \psi^2(x, \lambda) \int_0^x \psi^{-2}(t, \lambda) dt. \tag{28}$$

We set

$$\Psi(x, \lambda) = \int_0^x \psi^{-2}(t, \lambda) dt, x > 0, \lambda \in \Omega_{r\sigma}, \varepsilon \leq \sigma < \pi, r \geq r_\varepsilon. \tag{29}$$

**Lemma 2.** *Let  $\varepsilon \leq \sigma < \pi, r \geq r_\varepsilon$ . Then*

$$\begin{aligned} \Psi(x, \lambda) \sim & \frac{1}{2} \exp\left(2 \int_0^x (q_a(t) - \lambda)^{1/2} dt\right) \\ & \times \left[ 1 - 2 \int_x^{+\infty} \left( \alpha + \frac{\alpha'}{2\sqrt{q-\lambda}} \right) dt + Q(x, \lambda) \right], \end{aligned} \tag{30}$$

where

$$\sup_{x \geq 0, \lambda \in \Omega_{r\sigma}} |Q(x, \lambda)(x + a)^2(q - \lambda)^{5/2}| < \infty. \tag{31}$$

**Proof.** Fix  $\varepsilon > 0$  and everywhere until the end of the proof of the lemma we assume that  $\varepsilon < \sigma < \pi, \lambda \in \Omega_{r\sigma}, x > 0$ . The function  $\Psi$  can be represented as

$$\Psi(x, \lambda) = \int_0^{x/2} \psi^{-2}(t, \lambda) dt + \int_{x/2}^x \psi^{-2}(t, \lambda) dt =: \Psi_1(x, \lambda) + \Psi_2(x, \lambda). \tag{32}$$

Since for all  $\sigma - \pi < \arg q_a(x) - \lambda < \pi - \sigma$  then  $\forall 0 \leq t < x/2$

$$\Re \left( \int_t^x \sqrt{q_a - \lambda} d\tau \right) > \Re \left( \int_{x/2}^x \sqrt{q_a - \lambda} d\tau \right)$$

$$> \sin(\sigma/2) \min \left\{ \sqrt{q_a(x) - \lambda}, \sqrt{q_a(x/2) - \lambda} \right\}. \tag{33}$$

Futher, since  $\forall t \in [x/2, x] \quad |q_a(x) - q_a(t)| < \log 2$  and  $|q_a(x) - \lambda| > r \sin \sigma$  then for all  $r > 2 \log 2 / \sin \sigma$

$$1/2 < \left| \frac{q_a(x) - \lambda}{q_a(t) - \lambda} \right| < 3/2 \quad \forall t \in [x/2, x], \lambda \in \Omega_{r\sigma}. \tag{34}$$

Hence, taking into account (33), we conclude

$$\Re \left( \int_t^x \sqrt{q_a - \lambda} d\tau \right) \geq C |q_a(x) - \lambda|^{1/2} x.$$

Therefore, for all  $x \geq 0, \lambda \in \Omega_{r\sigma}$ , where  $r > 2 \log 2 / \sin \varepsilon$ ,

$$|\Psi_1(x, \lambda)| \leq C(r) \left| e^{2 \int_0^x (q_a(t) - \lambda)^{1/2} dt} (x + a)^{-2} (q(x) - \lambda)^{-5/2} \right|, \tag{35}$$

$C(\varepsilon) > 0$  depends only on  $\varepsilon$ .

Next, we substitute the expression (21) into the formula for  $\Psi_2$  and integrate the latter in parts. Cosequently, according to (34), we obtain (30), (31).

**Lemma 3.** *Let  $0 < \varepsilon \leq \sigma < \pi$ . Then there exists  $r_\varepsilon > 0$  so that the representation is true*

$$G(x, x, \lambda) = \frac{1}{2} (q_a(x) - \lambda)^{-1/2} \left( 1 - e^{-2 \int_0^x \sqrt{q_a(t) - \lambda} dt} + g(x, \lambda) \right), \tag{36}$$

$$|g(x, \lambda)| \leq C(\varepsilon) (x + a)^{-2} |q(x) - \lambda|^{-5/2} \quad \forall x \geq 0, \lambda \in \Omega_{r\sigma}, \tag{37}$$

where  $r > r_\varepsilon, C(\varepsilon) > 0$  depends only on  $\varepsilon$ .

**Proof.** It follows from the formula (28) and Lemmas 2, 3.

### 3.2. Formula for the kernel of the operator $e^{-tL_a}$

Let  $0 < \varepsilon \leq \sigma < \pi, r > r_\varepsilon$ , where  $r_\varepsilon$  is the constant appearing in the formulation of the Lemma 3,  $\gamma_{r\sigma}$  is the boundary of the region  $\Omega_{r\sigma}$  with the counterclockwise direction<sup>1</sup>.

**Lemma 4.** *For any  $t > 0$  the representation*

$$e^{-tL_a} = \frac{1}{2\pi i} \int_{\gamma_{r\sigma}} e^{-t\lambda} (\lambda - L_a)^{-1} d\lambda \tag{38}$$

holds. Here the limit is understood in the sense of convergence in the uniform operator topology.

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<sup>1</sup>Since  $\lambda_1 > \log a$ , then bypassing  $\gamma_{r\sigma}$  the entire spectrum of  $L_a$  remains on the left.



**Proof.** Since  $e^{-tL_a}$  is a bounded operator, then

$$e^{-tL_a} = \lim_{n \rightarrow \infty} e^{-tL_a} P_n,$$

where  $P_n$  is orthogonal projector on the linear span of the first  $n$  eigenvectors of the operator  $L_a$  and the limit is understood in the sense of a convergence in the norm. We have (see, for example, [14, § XII.2])

$$P_n = \frac{1}{2\pi i} \int_{\Gamma_n} (\lambda - L_a)^{-1} d\lambda,$$

where  $\gamma_n$  is any contour covering the first  $n$  eigenvalues of the operator  $L_a$ , the integral is taken in the counterclockwise direction. Then

$$e^{-tL_a} P_n = \frac{1}{2\pi i} \int_{\gamma_n} e^{-tL_a} (\lambda - L_a)^{-1} d\lambda.$$

As  $\gamma_n$  we take a contour composed of  $l_n$  - parts of  $\gamma_{r\sigma}$  lying in the half-plane  $\Re \lambda < (\lambda_n + \lambda_{n+1})/2$ , and the segment  $l_n$  connecting the ends of  $\gamma_n$ .

Let

$$R_n = \int_{l_n} e^{-tL_a} (\lambda - L_a)^{-1} d\lambda.$$

Let us prove that for any  $\Re t > 1$

$$\|R_n\| \rightarrow 0, \quad n \rightarrow \infty. \tag{39}$$

We have

$$\|(\lambda - L_a)^{-1}\| \leq \frac{1}{\text{dist}(\lambda, \sigma(L_a))}. \tag{40}$$

Hence, for all  $\lambda$  from  $[N_-, N_+]$ , where  $N_{\pm} = (\lambda_n + \lambda_{n+1})/2 \pm i$ , we will have  $\|(\lambda - L_a)^{-1}\| \leq 2/(\lambda_{n+1} - \lambda_n)$ . According to (1) - (2)  $\lambda_{n+1} - \lambda_n \sim n^{-1}$ ,  $n \rightarrow \infty$ . Hence, for the operator

$$r_n = \int_{[N_-, N_+]} e^{-tL_a} (\lambda - L_a)^{-1} d\lambda$$

we get

$$\|r_n\| = O(n^{-t+1}), \quad n \rightarrow \infty. \tag{41}$$

Further, since according to (40)  $\|(\lambda - L_a)^{-1}\| < 1$  on  $l_n \setminus [N_-, N_+]$ , then

$$\|R_n - r_n\| = O(e^{-tn} n), \quad n \rightarrow \infty. \tag{42}$$

Now (39) directly follows from (41) and (42).

**3.3. Asymptotics of  $\text{tr}(e^{-tL_a})$**

**Lemma 5.** For  $\Re t > 1$

$$\text{tr}(e^{-tL_a}) = -\frac{1}{2\pi i} \int_0^\infty dx \int_{\gamma_{r\sigma}} e^{-t\lambda} G(x, x, \lambda) d\lambda.$$

**Proof.** By Lemma 4

$$(e^{-tL_a} f)(x) = -\frac{1}{2\pi i} \int_\Gamma d\lambda e^{-t\lambda} \int_0^\infty G(x, y, \lambda) f(y) dy. \tag{43}$$

According to formulas (19), (27) and Lemmas 1 and 2

$$|G(x, y, \lambda)| \leq C |(q_a(x) - \lambda)^{-1/4} (q_a(y) - \lambda)^{-1/4} \exp\left(-\left|\Re \int_x^y \sqrt{q_a(t) - \lambda} dt\right|\right)|, \tag{44}$$

where  $C = \text{const} > 0$ . Further, for any fix  $x > 0$  and all  $f \in L^2(0, +\infty)$

$$\begin{aligned} |(e^{-tL_a} f)(x)| &\leq \int_\Gamma |d\lambda| |e^{-t\lambda}| \int_0^\infty |G(x, y, \lambda) f| dy \\ &\leq \frac{1}{2\pi} \int_\Gamma |d\lambda| |e^{-t\lambda}| \sqrt{\int_0^\infty |G(x, y, \lambda)|^2 dy} \cdot \|f\| \end{aligned}$$

Hence, taking into account the estimate (44), we will have

$$|e^{-tL_a} f(x)| \leq C \int_{\gamma_{r\sigma}} |e^{-t\lambda}| |q_a(x) - \lambda|^{-1/4} |d\lambda| < \infty.$$

Therefore, we can apply the Fubini theorem to the right side of (43). As a result, we get

$$e^{-tL_a} f = - \int_0^\infty \left( \frac{1}{2\pi i} \int_{\gamma_{r\sigma}} d\lambda e^{-t\lambda} G(x, y, \lambda) \right) f dy.$$

Consequently, the kernel of the operator  $e^{-tL_a}$  has the form

$$H(x, y, t) = -\frac{1}{2\pi i} \int_{\gamma_{r\sigma}} d\lambda e^{-t\lambda} G(x, y, \lambda).$$

According to formulas (1) – (2) for  $\Re t > 1$  the operator  $e^{-tL_a}$  is a trace-class one, therefore

$$\text{tr}(e^{-tL_a}) = \int_0^\infty H(x, x, t) dx = -\frac{1}{2\pi i} \int_0^\infty dx \int_{\gamma_{r\sigma}} d\lambda e^{-t\lambda} G(x, x, \lambda).$$

**Theorem 2.** For the  $\theta$ -functions of the operator  $L_a$  as  $t \rightarrow +0$  the following decomposition is true

$$\Theta_a(t) \sim -\frac{a}{2\sqrt{\pi}} t^{-1/2} - 1/4 + \theta_0 t^{1/2} + \frac{\log a}{4} t + O(t^{3/2}). \tag{45}$$

Here  $\theta_0$  is defined in the same way as in (6).

**Proof.** The statement of the theorem follows directly from the Lemmas 3 and 5.

**Corollary 3.** *The constant  $k_0$  in formulas (1) – (2) and (6) is equal to 0.*

#### 4. Regularized trace of the operator $L_a$

The sum

$$\sigma = \lambda_1 + \sum_{k=2}^{\infty} \left[ \lambda_k - \log(2\sqrt{\pi}k) - k^{-1} \left( \frac{a}{\pi} (\log k)^{1/2} - 1/4 + c_0 (\log k)^{-1/2} \right) \right], \quad (46)$$

where  $c_0 = \frac{a}{2\pi} (1 + \log(2\sqrt{\pi}/a))$ , we call the *regularized trace of the operator  $L_a$* .

**Theorem 4.** *The following formula is true*

$$\begin{aligned} \sigma &= \frac{1}{4} (\log(2\sqrt{\pi}) + \gamma - 1 + \log a) - \frac{a}{\pi} \int_1^{\infty} \{x\} ((\log x)^{1/2} x^{-1})' dx \\ &\quad - \frac{a}{2\pi} \left( 1 + \log \frac{2\sqrt{\pi}}{a} \right) \int_1^{\infty} \{x\} ((\log x)^{-1/2} x^{-1})' dx. \end{aligned} \quad (47)$$

**Proof.** The formula (47) follows from the formulas (6) and (45).

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UDepartment of Mathematics and Information Technology, Bashkir State University, Zaki Validi street, 83, Ufa, 450076, Russia.

E-mail: [Ishkin62@mail.ru](mailto:Ishkin62@mail.ru)

Department of Mathematics and Information Technology, Bashkir State University, Zaki Validi street, 83, Ufa, 450076, Russia.

E-mail: [l.matem2012@yandex.ru](mailto:l.matem2012@yandex.ru)